

CHAPTER THREE

Poisson's and Laplace's Equation

Poisson's and Laplace's Equation:

Poisson's and Laplace's Equation provides a method where by the potential function V can be determined.

$$\nabla \cdot D = \rho \quad \text{Maxwell first equation}$$

$$\nabla \cdot E = \frac{\rho}{\epsilon} \quad \text{is the medium is homogenous and isotropic}$$

$$\nabla \cdot (-\nabla V) = \frac{\rho}{\epsilon}$$

$$\nabla^2 V = \frac{-\rho}{\epsilon} \quad \dots \dots \dots \quad \text{Poisson's eq.}$$

If the region of interest contains charge in a known distribution, Poisson's equation used to determine the function.

Very often the region is charge free that is i.e., ($\rho = 0$) $\rightarrow \nabla^2 V = 0$, Laplace's eq.

$$\left. \begin{aligned} \nabla^2 V &= \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} \\ \nabla V &= \frac{\partial V}{\partial x} \mathbf{i} + \frac{\partial V}{\partial y} \mathbf{j} + \frac{\partial V}{\partial z} \mathbf{k} \end{aligned} \right\} \text{Cartesian coordinate}$$

$$\left. \begin{aligned} \nabla^2 V &= \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial V}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 V}{\partial \phi^2} + \frac{\partial^2 V}{\partial z^2} \\ \nabla V &= \frac{\partial V}{\partial r} \mathbf{a}_r + \frac{1}{r} \frac{\partial V}{\partial \phi} \mathbf{a}_\phi + \frac{\partial V}{\partial z} \mathbf{a}_z \end{aligned} \right\} \text{cylindrical coordinate}$$

$$\left. \begin{aligned} \nabla^2 V &= \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V}{\partial \phi^2} \\ \nabla V &= \frac{\partial V}{\partial r} \mathbf{a}_r + \frac{1}{r} \frac{\partial V}{\partial \theta} \mathbf{a}_\theta + \frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi} \mathbf{a}_\phi \end{aligned} \right\} \text{spherical coordinate}$$

Cartesian sol. in one variable:

$$V = V(x, y, z)$$

if $V = V(x)$ or $V = V(y)$ or $V = V(z)$ then:

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} \quad \text{or} \quad \nabla^2 V = \frac{\partial^2 V}{\partial y^2} \quad \text{or} \quad \nabla^2 V = \frac{\partial^2 V}{\partial z^2}$$

If $V = V(x)$ only

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} = 0 \quad \rightarrow \quad \int \frac{d}{dx} \left(\frac{dV}{dx} \right) = \int 0 \quad \rightarrow \quad \frac{dV}{dx} = A \quad \equiv \text{in one variable}$$

$$dV = A dx \quad \rightarrow \quad \int dV = \int A dx \quad \rightarrow \quad \therefore V = Ax + B$$

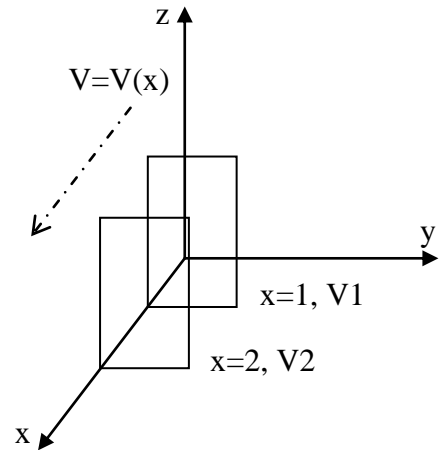
For a parallel conductor

Let: $V = V_1$ at $x = x_1$, $V = V_2$ at $x = x_2$

$$\left. \begin{aligned} V_1 &= Ax_1 + B \\ V_2 &= Ax_2 + B \end{aligned} \right\} \text{ subtract them } \rightarrow A = \frac{V_1 - V_2}{x_1 - x_2}$$

$$B = V_1 - Ax_1 = V_1 - \left(\frac{V_1 - V_2}{x_1 - x_2} \right) x_1 = \frac{V_2 x_1 - V_1 x_2}{x_1 - x_2}$$

$$V = Ax + B = \left(\frac{V_1 - V_2}{x_1 - x_2} \right) x + \frac{V_2 x_1 - V_1 x_2}{x_1 - x_2}$$



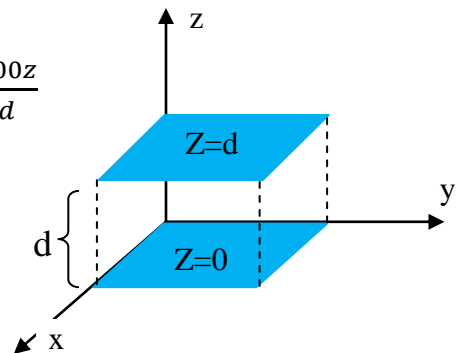
$$V = \frac{V_1(x - x_2) - V_2(x - x_1)}{x_1 - x_2} \quad \dots \dots \dots \quad \text{General solution}$$

Example: Consider the parallel conductors where $V=0$ at $Z=0$, $V=100$ volt at $Z=d$.

$$\text{Sol} \setminus V = \frac{V_1(z-z_2) - V_2(z-z_1)}{z_1 - z_2} = \frac{V(z-d) - 100(z-0)}{0-d} = \frac{100z}{d}$$

$$E = -\nabla V = -\frac{dV}{dz} a_z = -\frac{d}{dz} \left(\frac{100z}{d} \right) = \frac{100}{d}$$

$$D = \epsilon_0 E = \frac{100}{d} \epsilon_0$$



Ex.: Two Parallel conducting planes in free space are at $y=0$ and $y=2$ cm, and the zero voltage reference is at $y=1$ cm, if $D = 2.53j \frac{nC}{m^2}$ between the two conductors. Determine the conductor voltage.

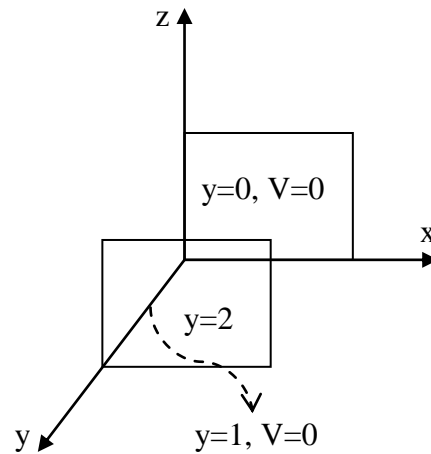
$$\text{Sol.} \quad \nabla^2 V = \frac{\partial^2 V}{\partial y^2} = 0$$

$$\int \frac{d}{dy} \left(\frac{dV}{dy} \right) = \int 0 \rightarrow V = Ay + B$$

$$E = -\nabla V \rightarrow E = -Aj \dots \dots \dots 1$$

$$D = \epsilon_0 E \rightarrow E = \frac{D}{\epsilon_0} \text{ in free space}$$

$$E = \frac{253 \times 10^{-9}}{8.8542 \times 10^{-12}} = 2.86 \times 10^4 \frac{V}{m}$$



From eq. 1 we get value of A

$$A = -2.86 \times 10^4 \frac{V}{m}, \quad \therefore V = -2.86 \times 10^4 y + B$$

Boundary condition at $y=1$ cm

$$0 = -2.86 \times 10^4 \times 10^{-2} + B \quad \therefore B = 286$$

$$V = -2.86 \times 10^4 y + 286$$

at $y=0$

$$V_1 = -2.86 \times 10^4 \times 0 + 286 \rightarrow V_1 = 286 \text{ volt}$$

at $y=2$ cm

$$V_2 = -2.86 \times 10^4 \times 2 \times 10^{-2} + 286 \rightarrow V_2 = -286 \text{ volt}$$

Cylindrical solution in one variable:

$$V = V(r, \varphi, z)$$

$$V = V(r) \quad \text{or} \quad V = V(\varphi) \quad \text{or} \quad V = V(z)$$

1- For $V = V(r)$ only

$$\nabla^2 V = \frac{1}{r} \frac{d}{dr} \left(r \frac{dV}{dr} \right) = 0$$

$$\int \frac{d}{dr} \left(r \frac{dV}{dr} \right) = 0 \quad \rightarrow \quad r \frac{dV}{dr} = A$$

$$\int dV = A \int \frac{dr}{r} \quad \rightarrow \quad V = A \ln r + B$$

As a Example:

For boundary condition $V = V_0$ at $r = a$
 $V = 0$ at $r = b$

Equipotential surface in concentric cylindrical.

$$\left. \begin{aligned} V_0 &= A \ln a + B \\ 0 &= A \ln b + B \end{aligned} \right\} \text{from them we get } A = \frac{V_0}{\ln a - \ln b} = \frac{V_0}{\ln \frac{a}{b}}$$

$$B = -A \ln b = \frac{-\ln b}{\ln \frac{a}{b}} V_0$$

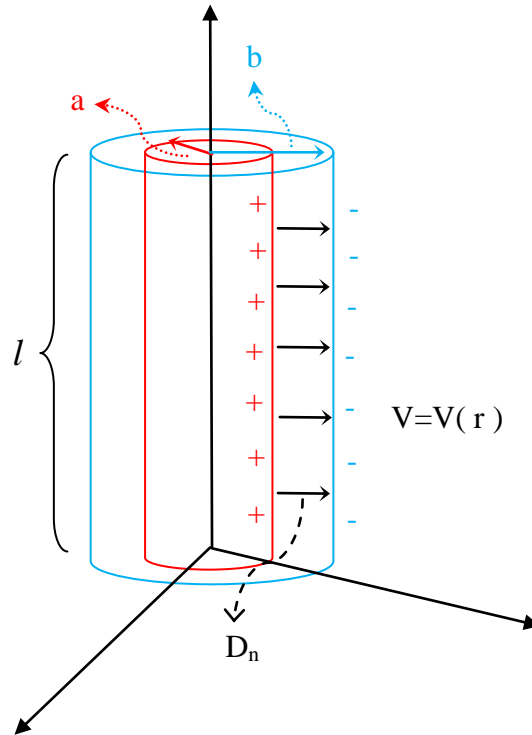
$$\therefore V = \frac{V_0}{\ln \frac{a}{b}} \ln r - \frac{\ln b}{\ln \frac{a}{b}} V_0$$

$$V = V_0 \frac{\ln \frac{r}{b}}{\ln \frac{a}{b}}$$

$$E = -\nabla V = -\frac{dV}{dr} a_r$$

$$= -\left[\frac{V_0}{\ln \frac{a}{b}} \frac{d}{dr} (\ln r - \ln b) \right] = \frac{V_0}{r \ln \frac{b}{a}}$$

$$D = \epsilon_0 E = \frac{\epsilon_0 V_0}{r \ln \frac{b}{a}}$$



At $r=a$

$$D_n = \frac{\epsilon_0 V_0}{r \ln \frac{b}{a}} D_n = flux = \frac{\varphi}{A} = \rho_s = \frac{\varphi}{2\pi a l}$$

$$\frac{\varphi}{2\pi a l} = \frac{\epsilon_0 V_0}{a \ln \frac{b}{a}} \quad \therefore \varphi = \frac{\epsilon_0 V_0}{a \ln \frac{b}{a}} \times 2\pi a l = \frac{\epsilon_0 V_0}{\ln \frac{b}{a}} 2\pi l$$

$$C = \frac{\varphi}{V_0} = \frac{\epsilon_0}{\ln \frac{b}{a}} 2\pi l$$

Q/ Find V, E & D for the region between two concentric cylinders where $V=0$ at $r = 1\text{mm}$, & $V=150\text{volt}$ at $r=20\text{mm}$.

2- For $V = V(\varphi)$ only, and the boundary condition might be:

$$\begin{aligned} V &= 0 & \text{at } \varphi &= 0 \\ V &= V_0 & \text{at } \varphi &= \alpha \end{aligned}$$

$$\nabla^2 V = \frac{1}{r^2} \frac{d^2 V}{d\varphi^2} = 0 \quad \int \frac{d}{d\varphi} \left(\frac{dV}{d\varphi} \right) = \int 0 \quad , \quad \frac{dV}{d\varphi} = A$$

$$dV = A d\varphi \quad \quad V = A\varphi + B$$

$$\text{at } V = 0 \text{ , } \varphi = 0 \text{ , } \quad \text{and } V = V_0 \text{ , } \varphi = \alpha$$

$$\left. \begin{aligned} 0 &= 0 + B \\ V_0 &= A \cdot \alpha + B \end{aligned} \right\} \begin{aligned} B &= 0 \\ A &= \frac{V_0}{\alpha} \end{aligned}$$

$$\therefore V = \frac{V_0}{\alpha} \varphi + 0$$

$$\nabla V = \frac{\partial V}{\partial r} a_r + \frac{1}{r} \frac{\partial V}{\partial \varphi} a_\varphi + \frac{\partial V}{\partial z} a_z$$

$$E = -\nabla V = -\frac{1}{r} \frac{\partial V}{\partial \varphi} a_\varphi = -\frac{1}{r} \frac{\partial}{\partial \varphi} \left(\frac{V_0}{\alpha} \varphi \right) a_\varphi = -\frac{1}{r} \frac{V_0}{\alpha} a_\varphi$$

$$E = \frac{V_0}{r\alpha} (-a_\varphi)$$

$$D = \epsilon_0 E = \frac{\epsilon_0 V_0}{r} \frac{1}{\alpha} (-a_\varphi)$$

Example: In cylindrical coordinate there are two plane charged ($\varphi = \text{constant}$), then planes are insulated along z-axis. Find the expression for D between the planes where: $V = 100v$ at $\varphi = \alpha$ and $V = 0$ at $\varphi = 0$

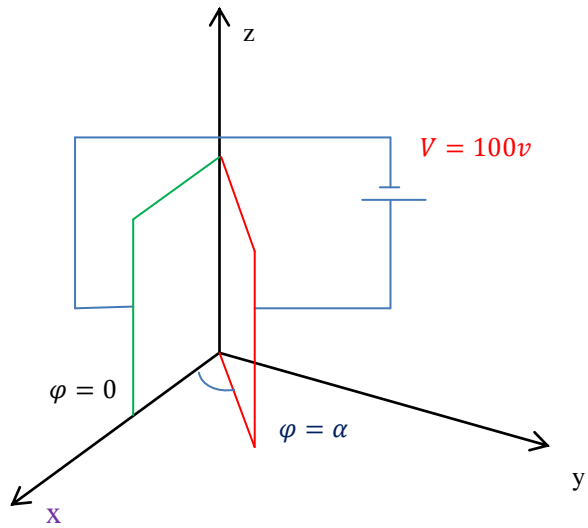
Sol/ $V = A\varphi + B$

$$\left. \begin{aligned} 100 &= A\alpha + B \\ 0 &= 0 + B \end{aligned} \right\} \therefore \begin{aligned} B &= 0 \\ A &= \frac{100}{\alpha} \end{aligned}$$

$$V = \frac{100\varphi}{\alpha}$$

$$E = -\nabla V = -\frac{1}{r} \frac{dV}{d\varphi} a_\varphi = \frac{-100}{r\alpha} a_\varphi$$

$$D = \frac{-100\epsilon_0}{r\alpha} a_\varphi$$



The equipotential surfaces are semicircular radial planes, $V =$

3- For $V=V(z)$ only

$$\nabla^2 V = \frac{d^2 V}{dz^2} = 0 \quad \rightarrow \quad V = Az + B$$

Spherical solution in one variable

$$V = V(r, \theta, \varphi)$$

1- For $V=V(r)$ only

$$\nabla^2 V = \frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{dV}{dr} \right) = 0$$

$$r^2 \frac{dV}{dr} = A \quad \rightarrow \quad dV = A \frac{dr}{r^2}$$

$$V = -\frac{A}{r} + B$$

Example: In spherical coordinates $V = 0$ at $r = 0.1m$ and $V = 100v$ at $r = 2m$.
Find E and D.

Sol./

$$V = -\frac{A}{r} + B$$

$$\left. \begin{array}{l} 0 = -\frac{A}{0.1} + B \\ 100 = -\frac{A}{2} + B \end{array} \right\} 100 = -\frac{A}{2} + \frac{A}{0.1} = A \left(\frac{20-1}{2} \right) \rightarrow A = 9.5$$

$$B = \frac{A}{0.1} = \frac{9.5}{0.1} = 95$$

$$\therefore V = \frac{-9.5}{r} + 95$$

2- For $V=V(\theta)$ only

$$\nabla^2 V = \frac{1}{r^2 \sin \theta} \frac{d}{d\theta} \left(\sin \theta \frac{dV}{d\theta} \right) = 0$$

$$\int \frac{d}{d\theta} \left(\sin \theta \frac{dV}{d\theta} \right) = \int 0$$

$$\sin \theta \frac{dV}{d\theta} = A \quad \rightarrow \quad \int dV = A \int \frac{d\theta}{\sin \theta}$$

$$V = A \ln \left(\tan \frac{\theta}{2} \right) + B$$

Example: For the region between the two co-axial cones, the cones vertex are insulates at $[r=0]$. Find the potential function for $V = V_1$ at $\theta = \theta_1$ and $V = 0$ at $\theta = \theta_2$.

Sol./

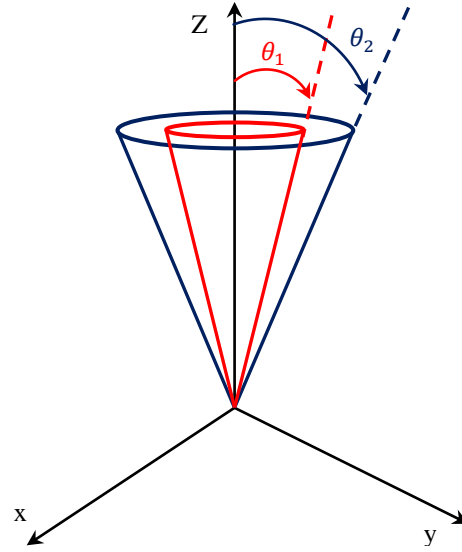
$$V = A \ln \left(\tan \frac{\theta}{2} \right) + B$$

$$\left. \begin{aligned} V_1 &= A \ln \left(\tan \frac{\theta_1}{2} \right) + B \\ 0 &= A \ln \left(\tan \frac{\theta_2}{2} \right) + B \end{aligned} \right\}$$

$$V_1 = A \left(\ln \left(\tan \frac{\theta_1}{2} \right) - \ln \left(\tan \frac{\theta_2}{2} \right) \right)$$

$$A = \frac{V_1}{\left(\ln \left(\tan \frac{\theta_1}{2} \right) - \ln \left(\tan \frac{\theta_2}{2} \right) \right)}$$

$$B = \frac{-V_1 \ln \left(\tan \frac{\theta_2}{2} \right)}{\left(\ln \left(\tan \frac{\theta_1}{2} \right) - \ln \left(\tan \frac{\theta_2}{2} \right) \right)}$$



The equipotential surfaces in spherical coordinates at $V = V(\theta)$ are coaxial cones

Then from above equations we get:

$$V = \frac{V_1 \ln \left(\tan \frac{\theta}{2} \right)}{\left(\ln \left(\tan \frac{\theta_1}{2} \right) - \ln \left(\tan \frac{\theta_2}{2} \right) \right)} - \frac{V_1 \ln \left(\tan \frac{\theta_2}{2} \right)}{\left(\ln \left(\tan \frac{\theta_1}{2} \right) - \ln \left(\tan \frac{\theta_2}{2} \right) \right)}$$

$$V = V_1 \frac{\ln \left(\tan \frac{\theta}{2} \right) - \ln \left(\tan \frac{\theta_2}{2} \right)}{\left(\ln \left(\tan \frac{\theta_1}{2} \right) - \ln \left(\tan \frac{\theta_2}{2} \right) \right)}$$

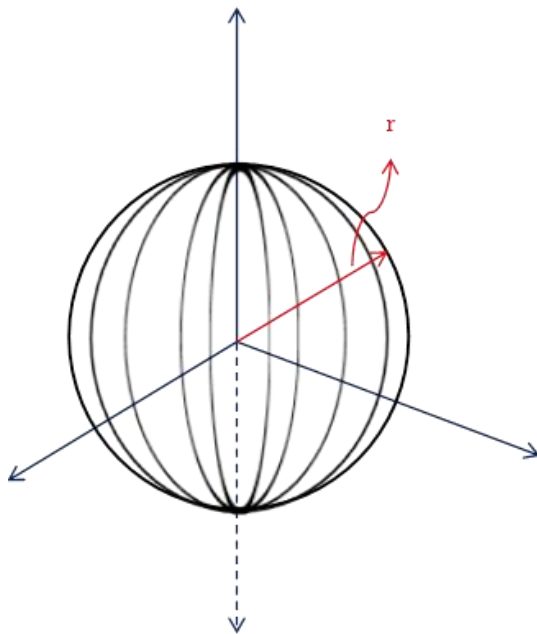
3- For $V = V(\varphi)$ only

$$\nabla^2 V = \frac{1}{r^2 \sin \theta} \frac{d^2 V}{d\varphi^2} = 0$$

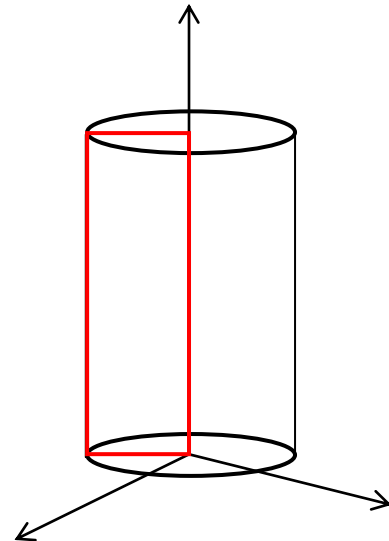
$$\frac{d^2 V}{d\varphi^2} = 0 \quad \rightarrow \quad \int \frac{d}{d\varphi} \left(\frac{dV}{d\varphi} \right) = \int 0 \quad \rightarrow \quad \frac{dV}{d\varphi} = A$$

$$\int dV = \int A d\varphi$$

$$V = A\varphi + B$$



*The equipotential surfaces
are semicircular radial planes,
 $V = V(\varphi)$*



*In cylindrical coordinate,
 $V = V(\varphi)$
The equipotential surfaces
are Rectangular radial
plane*

Solution of Poisson's eq.:

$$\nabla^2 V = \frac{-\rho}{\epsilon}$$

Example: The region between two concentric cylindrical contains a uniform charge density (ρ).

Sol./ For concentric cylindrical equipotential surfaces. $V = V(r)$

$$\nabla^2 V = \frac{1}{r} \frac{d}{dr} \left(r \frac{dV}{dr} \right) = \frac{-\rho}{\epsilon}$$

$$\int d \left(r \frac{dV}{dr} \right) = \int \frac{-\rho}{\epsilon} r dr$$

$$r \frac{dV}{dr} = \frac{-\rho}{2\epsilon} r^2 + A$$

$$\therefore \frac{dV}{dr} = \frac{-\rho r}{2\epsilon} + \frac{A}{r}$$

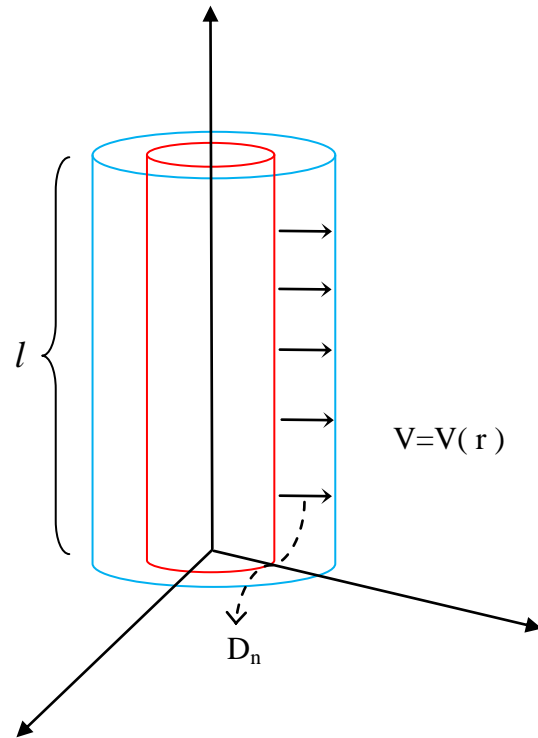
$$\int dV = \int \left(\frac{-\rho r}{2\epsilon} + \frac{A}{r} \right) dr$$

$$V = -\frac{\rho r^2}{4\epsilon} + A \ln r + B$$

$$E = -\nabla V = -\frac{dV}{dr} a_r$$

$$\therefore E = \frac{\rho r}{2\epsilon} - \frac{A}{r}$$

$$D = \epsilon E = \epsilon \left(\frac{\rho r}{2\epsilon} - \frac{A}{r} \right) = \frac{\rho r}{2} - \frac{\epsilon A}{r}$$



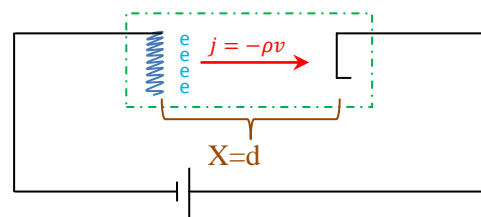
Example: A parallel plate Diode, operating condition x-axis, during the prose's space charge build up in the region between the electrons. Find the un expression for the potential

function between the plates if the current density $j = -\rho v$ where: $V = 0$ at $x = 0$
 $V = V_0$ at $x = d$

Sol./ charge density $\rho = -\frac{j}{v}$

$$eV = \frac{1}{2} m v^2 \quad \rightarrow \quad v = \sqrt{\frac{2eV}{m}}$$

$$\rho = \frac{-j}{\sqrt{\frac{2eV}{m}}}$$



$$\nabla^2 V = \frac{-\rho}{\varepsilon} \quad \rightarrow \quad \frac{d^2 V}{dx^2} = \frac{1}{\varepsilon} \frac{+j}{\sqrt{\frac{2eV}{m}}} = \frac{j}{\varepsilon} \left(\frac{2e}{m}\right)^{\frac{-1}{2}} V^{\frac{-1}{2}}$$

$$\frac{d^2 V}{dx^2} = k V^{\frac{-1}{2}} \quad \text{where } k = \frac{j}{\varepsilon} \left(\frac{2e}{m}\right)^{\frac{-1}{2}}$$

$$2 \left(\frac{dV}{dx}\right) \cdot \frac{d^2 V}{dx^2} = k V^{\frac{-1}{2}} \cdot 2 \left(\frac{dV}{dx}\right)$$

$$\int \frac{d}{dx} \left(\frac{dV}{dx}\right)^2 dx = \int 2k \left(\frac{dV}{dx}\right) V^{\frac{-1}{2}} dx$$

$$\left(\frac{dV}{dx}\right)^2 = 4kV^{\frac{+1}{2}} \quad \rightarrow \quad \frac{dV}{dx} = 2\sqrt{kV^{\frac{+1}{2}}}$$

$$\int V^{-\frac{1}{4}} dV = 2\sqrt{k} dx$$

$$\frac{4}{3} V^{\frac{3}{4}} = 2\sqrt{k} x + A$$

First boundary condition: $V = 0$ at $x = 0$

$$\frac{4}{3} (0)^{\frac{3}{4}} = 2\sqrt{k} \cdot 0 + A \quad \rightarrow \quad A = 0$$

$$\therefore \frac{4}{3} V^{\frac{3}{4}} = 2\sqrt{k} x \quad \dots \dots \dots 1$$

Second boundary condition: $V = V_0$ at $x = d$

$$\frac{4}{3} (V_0)^{\frac{3}{4}} = 2\sqrt{k} d \quad \dots \dots \dots 2$$

From eq. 1 & 2

$$\frac{\frac{4}{3} V^{\frac{3}{4}}}{\frac{4}{3} (V_0)^{\frac{3}{4}}} = \frac{2\sqrt{k} x}{2\sqrt{k} d} \quad \rightarrow \quad \frac{V^{\frac{3}{4}}}{(V_0)^{\frac{3}{4}}} = \frac{x}{d}$$

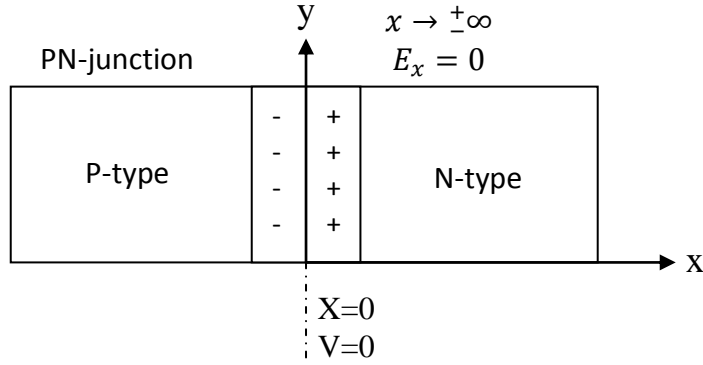
$$\therefore V = V_0 \left(\frac{x}{d}\right)^{\frac{4}{3}}$$

Q1/

A p-n junction between two halves of a semiconductor bar extended in the x-direction assume that the region for $x < 0$ is p-type and that the region for $x > 0$ is n-type. A charge distribution of this form may be approximated by the expression

$$\left(\rho = 2\rho_0 \operatorname{sech} \frac{x}{a} \tanh \frac{x}{a} \right), \text{ with}$$

$\rho_{max} = \rho_0$ where $\rho_{max} \equiv$ maximum charge density. Find V.



$$\nabla^2 V = -\frac{\rho_0}{\epsilon}$$

$$\nabla^2 V = \frac{2\rho_0 \operatorname{sech} \frac{x}{a} \tanh \frac{x}{a}}{\epsilon}$$

$$\frac{dV}{dx} = \frac{2a\rho_0 \operatorname{sech} \frac{x}{a}}{\epsilon} + C_1$$

$$E_x = \frac{2a\rho_0 \operatorname{sech} \frac{x}{a}}{\epsilon} + C_1$$

$$\text{at } x \rightarrow \pm\infty \quad E_x = 0, \quad \therefore C_1 = 0$$

$$E_x = \frac{dV}{dx} = \frac{2a\rho_0 \operatorname{sech} \frac{x}{a}}{\epsilon}$$

Integrating again:

$$V = \frac{4a^2\rho_0}{\epsilon} \tan^{-1} e^{\frac{x}{a}} + C_2$$

Let us arbitrarily select our zero reference of potential at the center of the junction, $= 0$;

$$0 = \frac{4a^2\rho_0}{\epsilon} \frac{\pi}{4} + C_2 \quad \therefore C_2 = -\frac{4a^2\rho_0}{\epsilon} \frac{\pi}{4}$$

$$V = \frac{4a^2\rho_0}{\epsilon} \left(\tan^{-1} e^{\frac{x}{a}} - \frac{\pi}{4} \right)$$

Q2/ Find E in the region between the two cones, where $\begin{cases} V = V_1 & \text{at } \theta_1 = 20^\circ \\ V = 0 & \text{at } \theta_1 = 160^\circ \end{cases}$.

Q3/ coaxial conducting cylindrical are located at $r=4\text{m}$ and 15cm , the value of E is

$\left(E = 20a_r \frac{kV}{m} \right)$ at $r = 6\text{cm}$, and the potential of the more positive conductors is 200V .

Find a- potential difference between two conductors.

b- the capacitance of the system if $\epsilon_r = 2.7$ for the medium between the two cylinders