Introduction to 3D Elasticity

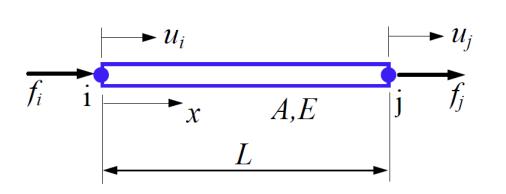
Summary:

- 1D elasticity (Bar Element)
- 3D elasticity problem
 - Governing differential equation
 - Strain-displacement relationship
 - Stress-strain relationship
- Special cases

2D (plane stress, plane strain)

Axisymmetric body with axisymmetric loading

1D Elasticity (Bar Element)



L = length

A = cross-sectional area

E = elastic modulus

u = u(x) displacement

 $\varepsilon = \varepsilon(x)$ strain

 $\sigma = \sigma(x)$ stress

Strain-displacement relation:

$$\varepsilon = \frac{du}{dx}$$

Stress-strain relation:

$$\sigma = E\varepsilon$$

Assuming that the displacement u is varying linearly along the axis of the bar, i.e.,

$$u(x) = \left(1 - \frac{x}{L}\right)u_i + \frac{x}{L}u_j$$

$$f_i \quad i \qquad x \quad A,E \quad j \quad f_j$$

we have

$$\varepsilon = \frac{u_j - u_i}{L} = \frac{\Delta}{L}$$
 ($\Delta = \text{ elongation}$), $\sigma = E\varepsilon = \frac{E\Delta}{L}$

We also have

$$\sigma = \frac{F}{A} (F = \text{force in bar})$$

Thus

$$F = \frac{EA}{I}\Delta = k\Delta$$

The bar is acting like a spring in this case, and we conclude that element stiffness matrix is

$$\mathbf{k} = \begin{bmatrix} k & -k \\ -k & k \end{bmatrix} = \begin{bmatrix} \frac{EA}{L} & -\frac{EA}{L} \\ -\frac{EA}{L} & \frac{EA}{L} \end{bmatrix}$$

or

$$\mathbf{k} = \frac{EA}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix}$$

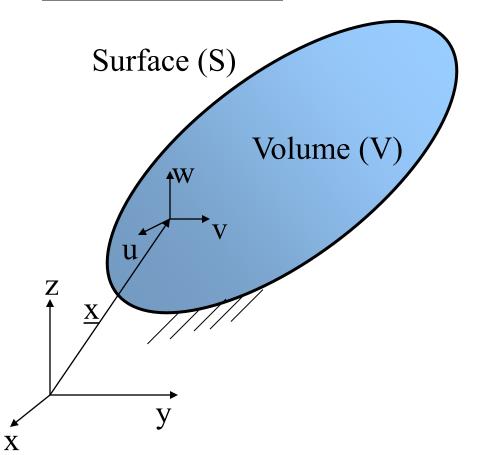
This can be verified by considering the equilibrium of the forces at the two nodes.

Element equilibrium equation is

$$\frac{EA}{L} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} u_i \\ u_j \end{Bmatrix} = \begin{Bmatrix} f_i \\ f_j \end{Bmatrix}$$

3D Elasticity

Problem definition



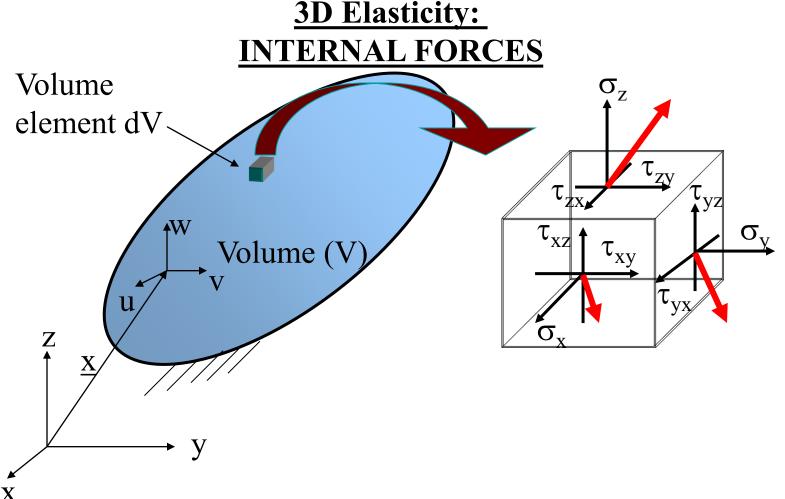
V: Volume of body

S: Total surface of the body

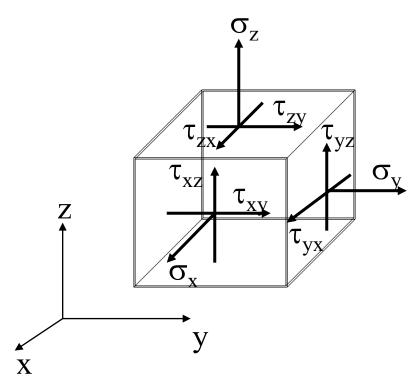
The deformation at point

$$\underline{x} = [x,y,z]^T$$
is given by the 3
components of its
displacement
 $\underline{u} = \begin{cases} u \\ v \end{cases}$

NOTE: $\underline{\mathbf{u}} = \underline{\mathbf{u}}(\mathbf{x}, \mathbf{y}, \mathbf{z})$, i.e., each displacement component is a function of position



If we take out a piece of material from the body, we will see that, due to the external forces applied to it, there are reaction forces (e.g., due to the loads applied to a truss structure, internal forces develop in each truss member). For the cube in the figure, the **internal reaction forces per unit area** (**red arrows**), on each surface, may be decomposed into three orthogonal components.



3D Elasticity

 σ_{x} , σ_{y} and σ_{z} are <u>normal stresses</u>. The rest 6 are the <u>shear stresses</u>

Convention τ_{xy} is the stress on the face perpendicular to the x-axis and points in the +ve y direction

Total of 9 stress components of which only 6 are $\tau_{yz} = \tau_{yx}$ $\tau_{yz} = \tau_{zy}$ independent since $\tau_{zx} = \tau_{xz}$

The stress vector is therefore
$$\underline{\sigma} = \begin{cases} \sigma_y \\ \sigma_z \\ \tau_{xy} \\ \tau_{yz} \\ \tau_{zx} \end{cases}$$

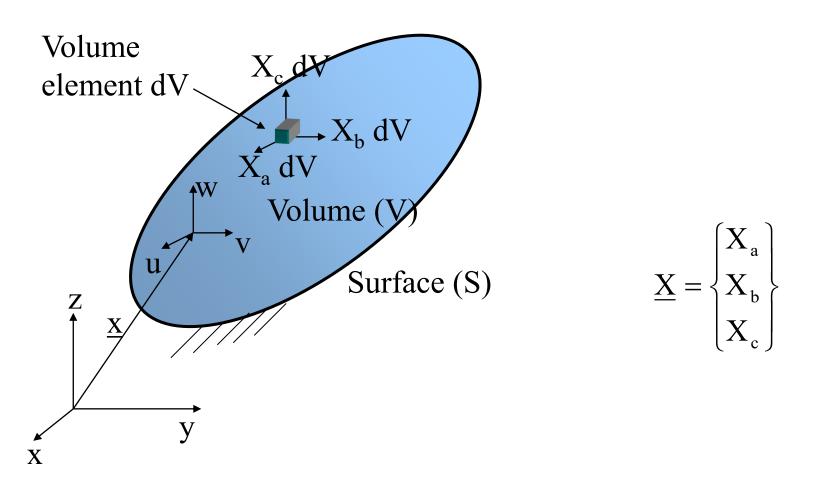
6 independent strain components

$$\underline{\varepsilon} = \begin{cases}
\varepsilon_y \\
\varepsilon_z \\
\gamma_{xy} \\
\gamma_{yz} \\
\gamma_{zx}
\end{cases}$$

BODY FORCE

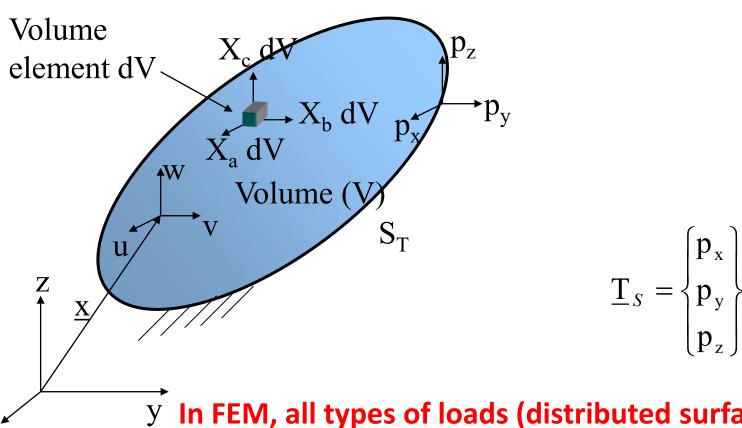
EXTERNAL FORCES ACTING ON THE BODY

1. Body force: distributed <u>force per unit volume</u> (e.g., weight, inertia, etc)



SURFACE TRACTION

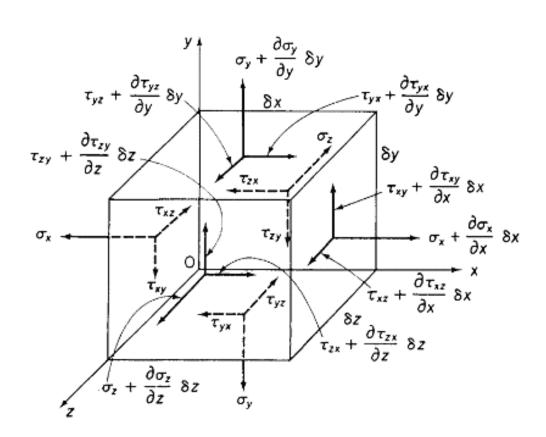
2. Surface traction (force per unit surface area) e.g., friction



y In FEM, all types of loads (distributed surface loads, body forces, concentrated forces and moments, etc.) are converted to point forces acting at the nodes.

Consider the equilibrium of a differential volume element to obtain the 3 equilibrium equations of elasticity

$$\frac{\partial \sigma_{x}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + \frac{\partial \tau_{xz}}{\partial z} + X_{a} = 0 \qquad \tau_{zy} + \frac{\partial \tau_{zy}}{\partial z} \delta z - \frac{\partial \tau_{xy}}{\partial z} + \frac{\partial \sigma_{y}}{\partial y} + \frac{\partial \tau_{yz}}{\partial z} + X_{b} = 0 \qquad \sigma_{x} + \frac{\partial \tau_{zy}}{\partial z} \delta z - \frac{\partial \tau_{xy}}{\partial z} + \frac{\partial \sigma_{yz}}{\partial z} + \frac{\partial \sigma_{z}}{\partial z} + X_{b} = 0$$



Compactly;

$$\underline{\partial}^T \underline{\sigma} + \underline{X} = \underline{0} \tag{1}$$

where

$$\frac{\partial}{\partial x} = \begin{bmatrix} \frac{\partial}{\partial x} & 0 & 0 \\ 0 & \frac{\partial}{\partial y} & 0 \\ 0 & 0 & \frac{\partial}{\partial z} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial z} & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} & 0 & \frac{\partial}{\partial x} \end{bmatrix}$$

Strain-displacement relationships in 3D elasticity problem:

$$\varepsilon_{x} = \frac{\partial u}{\partial x}$$

$$\varepsilon_{y} = \frac{\partial v}{\partial y}$$

$$\varepsilon_{z} = \frac{\partial w}{\partial z}$$

$$\gamma_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}$$

$$\gamma_{yz} = \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y}$$

$$\gamma_{zx} = \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x}$$

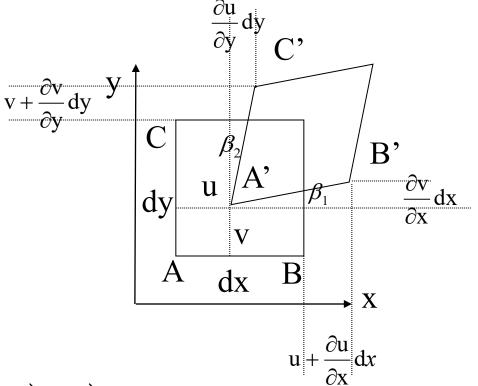
$$\underline{\varepsilon} = \underline{\partial} \ \underline{u}$$

$$\underline{\boldsymbol{\varepsilon}} = \begin{cases} \boldsymbol{\varepsilon}_{x} \\ \boldsymbol{\varepsilon}_{y} \\ \boldsymbol{\varepsilon}_{z} \\ \boldsymbol{\gamma}_{xy} \\ \boldsymbol{\gamma}_{yz} \\ \boldsymbol{\gamma}_{zx} \end{cases}$$

$$\frac{\partial}{\partial x} = \begin{bmatrix} \frac{\partial}{\partial x} & 0 & 0 \\ 0 & \frac{\partial}{\partial y} & 0 \\ 0 & 0 & \frac{\partial}{\partial z} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial z} & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} & 0 & \frac{\partial}{\partial x} \end{bmatrix}$$

$$\underline{\mathbf{u}} = \left\{ \begin{aligned} \mathbf{u} \\ \mathbf{v} \\ \mathbf{w} \end{aligned} \right\}$$

In 2D



$$\varepsilon_{x} = \frac{A'B' - AB}{AB} = \frac{\left(dx + \left(u + \frac{\partial u}{\partial x}dx\right) - u\right) - dx}{dx} = \frac{\partial u}{\partial x}$$

$$\varepsilon_{y} = \frac{A'C' - AC}{AC} = \frac{\left(dy + \left(v + \frac{\partial v}{\partial y}dy\right) - v\right) - dy}{dy} = \frac{\partial v}{\partial y}$$

$$\gamma_{xy} = \frac{\pi}{2} - \text{ angle } (C'A'B') = \beta_{1} + \beta_{2} \approx \tan\beta_{1} + \tan\beta_{2}$$

$$\approx \frac{\partial v}{\partial x} + \frac{\partial u}{\partial x}$$

Stress-Strain relationship in 3D elasticity problem:

Linear elastic material (Hooke's Law)

$$\underline{\sigma} = \underline{D} \, \underline{\varepsilon} \tag{3}$$

Linear elastic isotropic material

$$\underline{D} = \frac{E}{(1+\nu)(1-2\nu)} \begin{bmatrix} 1-\nu & \nu & \nu & 0 & 0 & 0\\ \nu & 1-\nu & \nu & 0 & 0 & 0\\ \nu & \nu & 1-\nu & 0 & 0 & 0\\ 0 & 0 & 0 & \frac{1-2\nu}{2} & 0 & 0\\ 0 & 0 & 0 & 0 & \frac{1-2\nu}{2} & 0\\ 0 & 0 & 0 & 0 & 0 & \frac{1-2\nu}{2} \end{bmatrix}$$

Plane Stress and Plane Strain

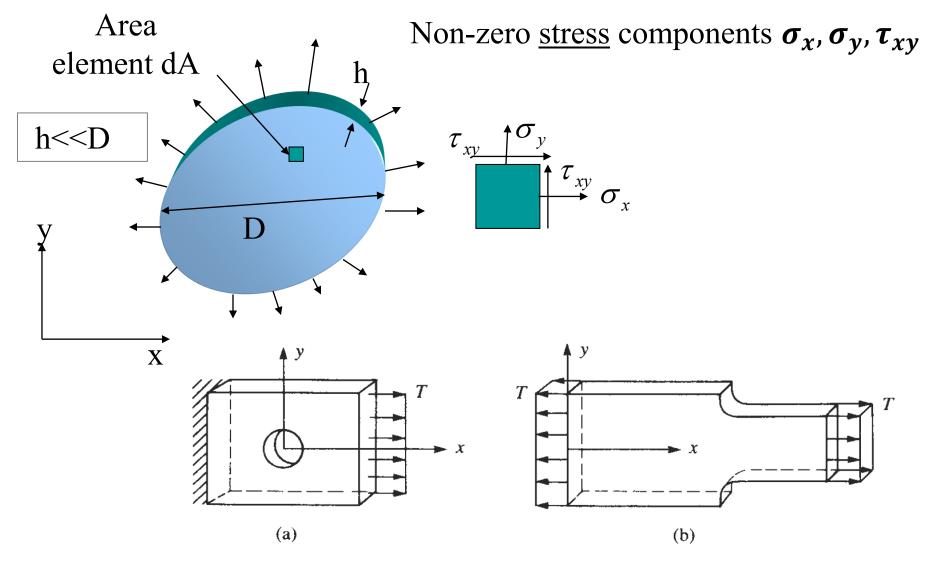
Plane stress is defined to be a state of stress in which the normal stress and the shear stresses directed perpendicular to the plane are assumed to be zero.

$$\sigma_z = \tau_{xz} = \tau_{yz} = 0$$

Plane strain is defined to be a state of strain in which the strain normal to the x-y plane ε_z and the shear strains γ_{xz} and γ_{yz} are assumed to be zero.

$$\varepsilon_z = \gamma_{xz} = \gamma_{yz} = 0$$

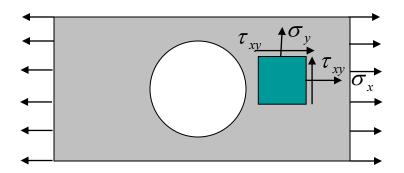
PLANE STRESS: Only the in-plane stress components are nonzero



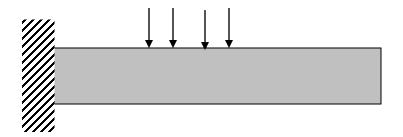
■ Figure 6–1 Plane stress problems: (a) plate with hole; (b) plate with fillet

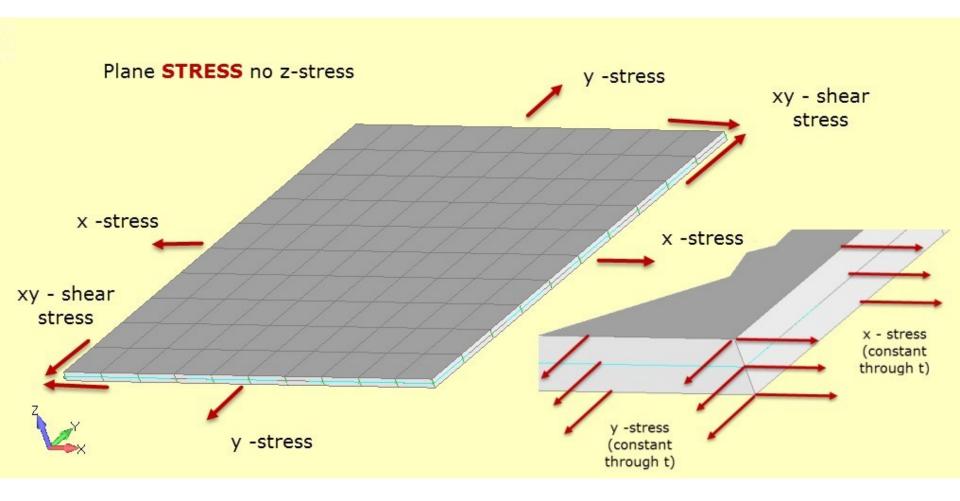
PLANE STRESS Examples:

1. Thin plate with a hole



2. Thin cantilever plate





PLANE STRESS

Non-zero <u>stresses</u>: $\sigma_x, \sigma_y, \tau_{xy}$

Non-zero <u>strains</u>: $\mathcal{E}_x, \mathcal{E}_y, \mathcal{E}_z, \gamma_{xy}$

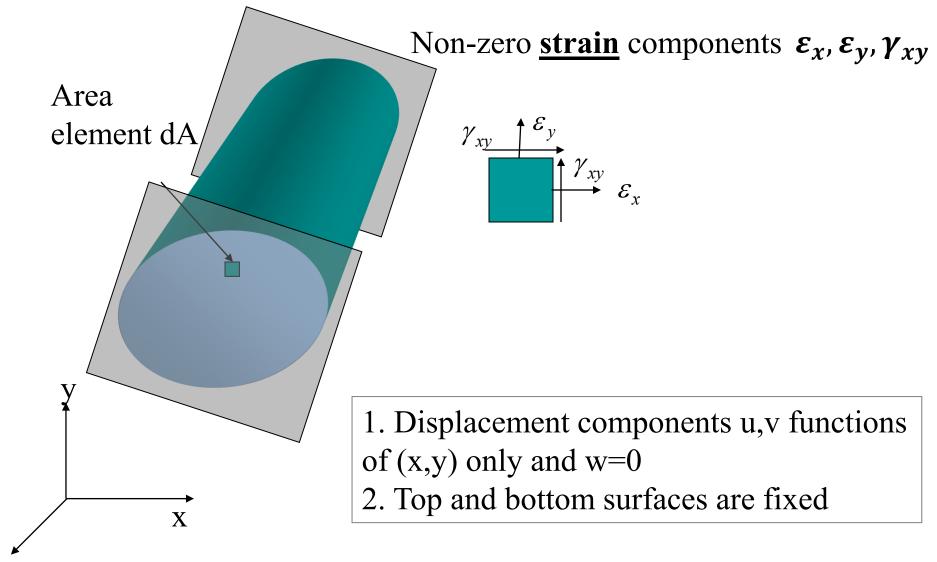
Isotropic linear elastic stress-strain law $\underline{\sigma} = \underline{D}$

$$\begin{cases}
\sigma_{x} \\
\sigma_{y} \\
\tau_{xy}
\end{cases} = \frac{E}{1 - \nu^{2}} \begin{bmatrix}
1 & \nu & 0 \\
\nu & 1 & 0 \\
0 & 0 & \frac{1 - \nu}{2}
\end{bmatrix} \begin{cases}
\varepsilon_{x} \\
\varepsilon_{y} \\
\gamma_{xy}
\end{cases} \qquad \varepsilon_{z} = -\frac{\nu}{1 - \nu} (\varepsilon_{x} + \varepsilon_{y})$$

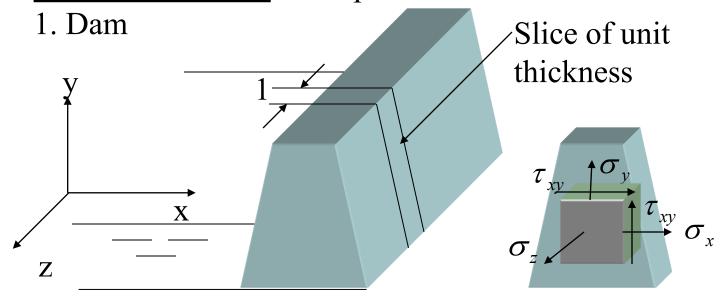
Hence, the <u>D</u> matrix for the <u>plane stress case</u> is

$$\underline{D} = \frac{E}{1 - v^2} \begin{bmatrix} 1 & v & 0 \\ v & 1 & 0 \\ 0 & 0 & \frac{1 - v}{2} \end{bmatrix}$$

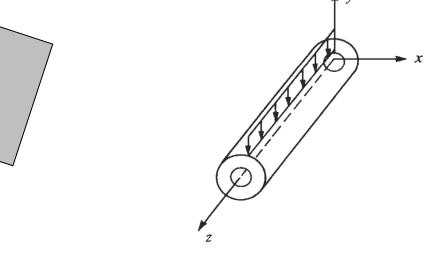
PLANE STRAIN: Only the in-plane strain components are non-zero



PLANE STRAIN Examples:



2. Long cylindrical pressure vessel subjected to internal/external pressure and constrained at the ends



PLANE STRAIN

Nonzero <u>stress</u>: $\sigma_x, \sigma_y, \sigma_z, \tau_{xy}$

Nonzero <u>strain</u> components: $\mathcal{E}_x, \mathcal{E}_y, \gamma_{xy}$

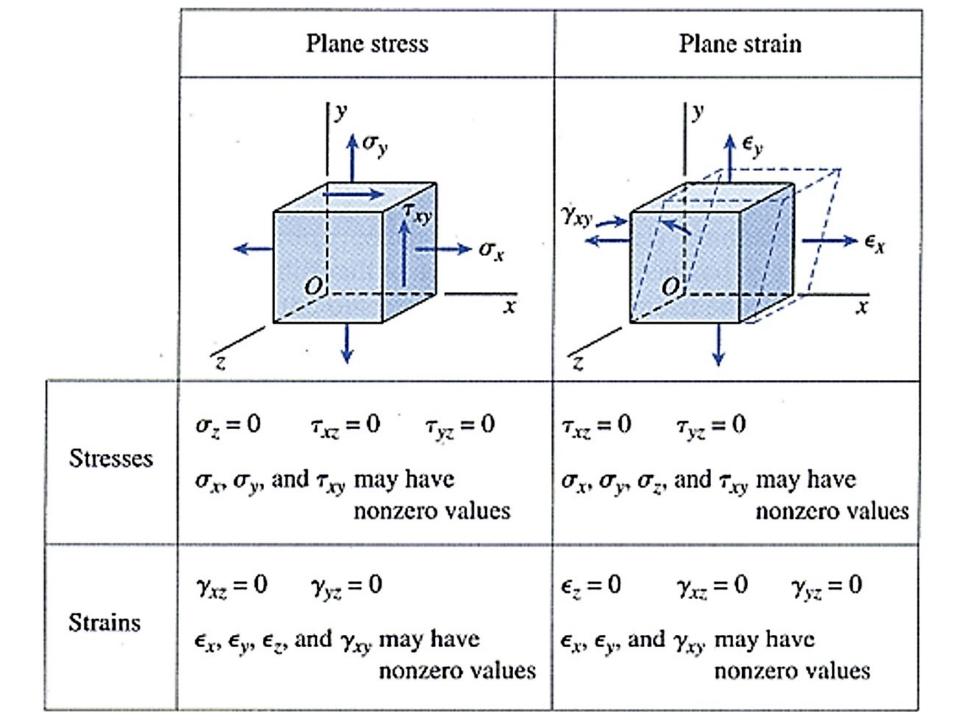
Isotropic linear elastic stress-strain law $|\underline{\sigma} = \underline{D} \underline{\varepsilon}|$

$$\underline{\sigma} = \underline{D} \, \underline{\varepsilon}$$

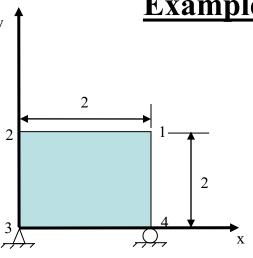
$$\begin{cases}
\sigma_{x} \\
\sigma_{y} \\
\tau_{xy}
\end{cases} = \frac{E}{(1+\nu)(1-2\nu)} \begin{bmatrix}
1-\nu & \nu & 0 \\
\nu & 1-\nu & 0 \\
0 & 0 & \frac{1-2\nu}{2}
\end{bmatrix} \begin{bmatrix}
\varepsilon_{x} \\
\varepsilon_{y} \\
\gamma_{xy}
\end{cases} \qquad \sigma_{z} = \nu \left(\sigma_{x} + \sigma_{y}\right)$$

Hence, the <u>D</u> matrix for the **plane strain case** is

$$\underline{D} = \frac{E}{(1+\nu)(1-2\nu)} \begin{vmatrix} 1-\nu & \nu & 0\\ \nu & 1-\nu & 0\\ 0 & 0 & \frac{1-2\nu}{2} \end{vmatrix}$$



Example problem



The square block is in **plane strain** and is subjected to the following strains

$$\varepsilon_x = 2xy$$

$$\varepsilon_y = 3xy^2$$

$$\gamma_{xy} = x^2 + y^2$$

Compute the displacement field (i.e., displacement components u(x,y) and v(x,y)) within the block

Solution

Recall from definition

$$\varepsilon_x = \frac{\partial u}{\partial x} = 2xy \qquad (1)$$

$$\varepsilon_y = \frac{\partial v}{\partial y} = 3xy^2 \qquad (2)$$

$$\gamma_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} = x^2 + y^3$$
 (3)

Integrating (1) and (2)

$$u(x, y) = x^2 y + C_1(y)$$
 (4)

$$v(x, y) = xy^3 + C_2(x)$$
 (5)

Arbitrary function of 'x'

Arbitrary function of 'y'

Plug expressions in (4) and (5) into equation (3)

$$\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} = x^2 + y^3 \quad (3)$$

$$\Rightarrow \frac{\partial \left[x^2 y + C_1(y)\right]}{\partial y} + \frac{\partial \left[xy^3 + C_2(x)\right]}{\partial x} = x^2 + y^3$$

$$\Rightarrow x^2 + \frac{\partial C_1(y)}{\partial y} + y^3 + \frac{\partial C_2(x)}{\partial x} = x^2 + y^3$$

$$\Rightarrow \frac{\partial C_1(y)}{\partial y} + \frac{\partial C_2(x)}{\partial x} = 0$$
Function of 'y'
Function of 'x'

Hence

$$\frac{\partial C_1(y)}{\partial y} = -\frac{\partial C_2(x)}{\partial x} = C \text{ (a constant)}$$

Integrate to obtain

$$C_1(y) = Cy + D_1$$
 D_1 and D_2 are two constants of $C_2(x) - Cx + D_2$ integration

Plug these back into equations (4) and (5)

(4)
$$u(x, y) = x^2y + Cy + D_1$$

(5)
$$v(x, y) = xy^3 - Cx + D_2$$

How to find C, D_1 and D_2 ?

Use the 3 **boundary conditions**

$$u(0,0) = 0$$

$$v(0,0) = 0$$

$$v(2,0) = 0$$

To obtain

$$C = 0$$

$$D_1 = 0$$

$$D_2 = 0$$

Hence the solution is

$$u(x, y) = x^{2}y$$
$$v(x, y) = xy^{3}$$

